

Weak and strong solutions to problems arising in dynamics of inviscid fluids

Eduard Feireisl

joint work with O.Kreml (Praha), Y.Sun (Nanjing), A.Vasseur (Austin),

Institute of Mathematics, Academy of Sciences of the Czech Republic, Prague

AMS/EMS/SPM 2015
Porto, 10 June - 13 June 2015

The research leading to these results has received funding from the European Research Council under the European Union's Seventh Framework Programme (FP7/2007-2013)/ ERC Grant Agreement 320078

Motivation

Compressible Euler system

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p(\varrho) = 0$$

Energy (entropy) inequality

$$\partial_t \left(\frac{1}{2} \varrho |\mathbf{u}|^2 + P(\varrho) \right) + \operatorname{div}_x \left[\left(\frac{1}{2} \varrho |\mathbf{u}|^2 + P(\varrho) \right) \mathbf{u} + p(\varrho) \mathbf{u} \right] \leq 0$$

$$P(\varrho) = \varrho \int_1^\varrho \frac{p(z)}{z^2} dz$$

Result of Chiodaroli, DeLellis, Kreml [2013]

There exist Lipschitz initial data such that the compressible Euler system admits infinitely many admissible (entropy) weak solutions.

Riemann problem

Riemann initial data

$$\varrho(0, x_1, \dots, x_N) = \begin{cases} \varrho_L & \text{if } x_1 \leq 0 \\ \varrho_R & \text{if } x_1 > 0 \end{cases}$$

$$u^1(0, x_1, \dots, x_N) = \begin{cases} u_L^1 & \text{if } x_1 \leq 0 \\ u_R^1 & \text{if } x_1 > 0 \end{cases}$$

$$u^k(0, x_1, \dots, x_N) = 0, \quad k = 2, \dots, N$$

Wild solutions

The wild solutions emanate from the 1D Riemann data but the velocity admits non-zero second component

Shock free solutions

Geometry, pressure

$$\Omega = (-a, a) \times \mathcal{T}^1 \text{ (periodic in } x_2)$$
$$p(0) = 0, \quad p'(r) > 0 \text{ for } r > 0, \quad p \text{ convex}$$

Theorem EF, O.Kreml [2014]

Let $\tilde{\varrho} = \tilde{\varrho}(x_1/t)$, $\tilde{\mathbf{u}} = [\tilde{u}^1(x_1/t), 0]$ be the self-similar solution to the Riemann problem consisting of rarefaction waves (locally Lipschitz for $t > 0$) and such that

$$\text{ess inf}_{(0,t) \times \mathcal{R}} \tilde{\varrho} > 0.$$

Let $[\varrho, \mathbf{u}]$ be a bounded admissible weak solution such that

$$\varrho \geq 0 \text{ a.a. in } (0, T) \times \Omega.$$

Then

$$\varrho \equiv \tilde{\varrho}, \quad \mathbf{u} \equiv \tilde{\mathbf{u}} \text{ in } (0, T) \times \Omega.$$

Method of relative energy

Relative energy

$$\mathcal{E}(\varrho, \mathbf{u} | r, \mathbf{U}) = \int_{\Omega} \left(\frac{1}{2} \varrho |\mathbf{u} - \mathbf{U}|^2 + P(\varrho) - P'(r)(\varrho - r) - P(r) \right) dx$$

Relative energy inequality

$$\begin{aligned} & \left[\mathcal{E}(\varrho, \mathbf{u} | \tilde{\varrho}, \tilde{\mathbf{u}}) \right]_{t=0}^{t=\tau} \\ & \leq - \int_0^{\tau} \int_{\Omega} \left[\varrho |u^1 - \tilde{u}^1|^2 + p(\varrho) - p'(\tilde{\varrho})(\varrho - \tilde{\varrho}) - p(\tilde{\varrho}) \right] \partial_{x_1} \tilde{u}^1 dx dt \\ & \quad + \text{"other terms"} \end{aligned}$$

Full Euler system

Mass conservation

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

Momentum balance

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x \varrho \vartheta = 0$$

Energy balance

$$\partial_t \left[\frac{1}{2} \varrho |\mathbf{u}|^2 + c_v \varrho \vartheta \right] + \operatorname{div}_x \left[\left(\frac{1}{2} \varrho |\mathbf{u}|^2 + c_v \varrho \vartheta + \varrho \vartheta \right) \mathbf{u} \right] = 0$$

Entropy inequality

$$\partial_t(\varrho s) + \operatorname{div}_x(\varrho s \mathbf{u}) \geq 0, \quad s = s(\varrho, \vartheta) \equiv \log \left(\frac{\vartheta^{c_v}}{\varrho} \right)$$

Riemann problem

Geometry

$\Omega = \mathbb{R}^1 \times \mathcal{T}^1$, where $\mathcal{T}^1 \equiv [0, 1]_{\{0,1\}}$ is the “flat” sphere

Initial data

$$\varrho(0, x_1, x_2) = R_0(x_1), \quad R_0 = \begin{cases} R_L & \text{for } x_1 \leq 0 \\ R_R & \text{for } x_1 > 0 \end{cases}$$

$$\vartheta(0, x_1, x_2) = \Theta_0(x_1), \quad \Theta_0 = \begin{cases} \Theta_L & \text{for } x_1 \leq 0 \\ \Theta_R & \text{for } x_1 > 0 \end{cases}$$

$$u^1(0, x_1, x_2) = U_0(x_1), \quad U_0 = \begin{cases} U_L & \text{for } x_1 \leq 0, \\ U_R & \text{for } x_1 > 0 \end{cases} \quad u^2(0, x_1, x_2) = 0.$$

Shock free Riemann solutions

Solution class

$$0 < \varrho \leq \bar{\varrho}, \quad 0 < \vartheta \leq \bar{\vartheta}, \quad |s(\varrho, \vartheta)| < \bar{s}, \quad |\mathbf{u}| < \bar{u}$$

Isentropic solutions

- the entropy S is *constant* in $[0, T] \times \Omega$
- $\Theta = R^{\frac{1}{c_v}} \exp\left(\frac{1}{c_v} S\right)$
- $R = R(t, x_1)$ and $U = U(t, x_1)$ represent a rarefaction wave solution of the 1-D *isentropic* system

$$\partial_t R + \partial_{x_1}(RU) = 0, \quad R[\partial_t U + U\partial_{x_1} U] + \exp\left(\frac{1}{c_v} S\right) \partial_{x_1} R^{\frac{c_v+1}{c_v}} = 0$$

Uniqueness

Theorem, EF, O.Kreml, A.Vasseur [2014]

Let $[\varrho, \vartheta, \mathbf{u}]$ be a weak solution of the Euler system in $(0, T) \times \Omega$ originating from the Riemann data. Suppose in addition that the Riemann data give rise to the shock-free solution $[R, \Theta, U]$ of the 1-D Riemann problem.

Then

$$\varrho = R, \vartheta = \Theta, \mathbf{u} = [U, 0] \text{ a.a. in } (0, T) \times \Omega$$

Relative energy

Relative energy (entropy) functional

$$\begin{aligned} & \mathcal{E}(\varrho, \vartheta, \mathbf{u} \mid \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \\ &= \int_{\Omega} \left[\frac{1}{2} \varrho |\mathbf{u} - \tilde{\mathbf{u}}|^2 + H_{\tilde{\vartheta}}(\varrho, \vartheta) - \frac{\partial H_{\tilde{\vartheta}}(\tilde{\varrho}, \tilde{\vartheta})}{\partial \varrho} (\varrho - \tilde{\varrho}) - H_{\tilde{\vartheta}}(\tilde{\varrho}, \tilde{\vartheta}) \right] dx \end{aligned}$$

Ballistic free energy

$$H_{\tilde{\vartheta}}(\varrho, \vartheta) = \varrho \left(c_v \vartheta - \tilde{\vartheta} s(\varrho, \vartheta) \right).$$

$$\varrho \mapsto H_{\tilde{\vartheta}}(\varrho, \tilde{\vartheta}) \text{ convex}$$

$$\vartheta \mapsto H_{\tilde{\vartheta}}(\varrho, \vartheta) \begin{cases} \text{decreasing for } \vartheta < \tilde{\vartheta} \\ \text{increasing for } \vartheta > \tilde{\vartheta} \end{cases}$$

Relative energy inequality, dissipative solutions

Relative energy inequality

$$\left[\mathcal{E}(\varrho, \vartheta, \mathbf{u} \mid \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \right]_{t=0}^{t=\tau} \leq \int_0^\tau \mathcal{R}(\varrho, \vartheta, \mathbf{u}, \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \, dt$$

Test functions

$$\tilde{\varrho} > 0, \tilde{\vartheta} > 0, \left\{ \begin{array}{l} \tilde{\varrho} = R_L, \tilde{\vartheta} = \Theta_L, \tilde{u}^1 = U_L, \tilde{u}^2 = 0 \text{ if } x_1 < -A, \\ \tilde{\varrho} = R_R, \tilde{\vartheta} = \Theta_R, \tilde{u}^1 = U_R, \tilde{u}^2 = 0 \text{ if } x_1 > A \end{array} \right\}$$

Remainder

Remainder in the relative energy inequality

$$\begin{aligned} & \mathcal{R}(\varrho, \vartheta, \mathbf{u}, \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \\ &= \int_{\Omega} \left[\varrho(\tilde{\mathbf{u}} - \mathbf{u}) \cdot \partial_t \tilde{\mathbf{u}} + \varrho(\tilde{\mathbf{u}} - \mathbf{u}) \otimes \mathbf{u} : \nabla_x \tilde{\mathbf{u}} + (\tilde{\varrho}\tilde{\vartheta} - \varrho\vartheta) \operatorname{div}_x \tilde{\mathbf{u}} \right] dx \\ & - \int_{\Omega} \left[\varrho \left(s(\varrho, \vartheta) - s(\tilde{\varrho}, \tilde{\vartheta}) \right) \partial_t \tilde{\vartheta} + \varrho \left(s(\varrho, \vartheta) - s(\tilde{\varrho}, \tilde{\vartheta}) \right) \mathbf{u} \cdot \nabla_x \tilde{\vartheta} \right] dx \\ & + \int_{\Omega} \left[\left(1 - \frac{\varrho}{\tilde{\varrho}} \right) \partial_t (\tilde{\varrho}\tilde{\vartheta}) + \left(\tilde{\mathbf{u}} - \frac{\varrho}{\tilde{\varrho}} \mathbf{u} \right) \cdot \nabla_x (\tilde{\varrho}\tilde{\vartheta}) \right] dx \end{aligned}$$

Robustness of 1D viscosity solutions

Navier-Stokes system

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p(\varrho) = \operatorname{div}_x \mathbb{S}(\nabla_x \mathbf{u}),$$

Pressure, viscous stress

$$p(\varrho) = a\varrho^\gamma, \quad a > 0, \quad \gamma > 1,$$

$$\mathbb{S}(\nabla_x \mathbf{u}) = \mu \left(\nabla_x \mathbf{u} + \nabla_x^t \mathbf{u} - \frac{2}{N} \operatorname{div}_x \mathbf{u} \mathbb{I} \right) + \eta \operatorname{div}_x \mathbf{u} \mathbb{I}, \quad \mu > 0, \quad \eta \geq 0.$$

1D problem

1D Navier-Stokes system

$$\partial_t R + \partial_y(RV) = 0,$$

$$\partial_t(RV) + \partial_y(RV^2) + \partial_y p(R) = \left[2\mu \left(1 - \frac{1}{N} \right) + \eta \right] \partial_{y,y}^2 V.$$

Stability of 1D solutions - hypotheses

Theorem EF, Y.Sun [2015]

$$\gamma > \frac{N}{2}, \quad q > \max\{2, \gamma'\}, \quad \frac{1}{\gamma} + \frac{1}{\gamma'} = 1 \text{ if } N = 2$$

$$q > \max\left\{3, \frac{6\gamma}{5\gamma - 6}\right\} \text{ if } N = 3$$

Let $[R, V]$ be a (strong) solution of the one-dimensional problem, with the initial data belonging to the class

$$R_0 \in W^{1,q}(0,1), \quad R_0 > 0, \quad V_0 \in W_0^{1,q}(0,1)$$

Let $[\varrho, \mathbf{u}]$ be a finite energy weak solution to the Navier-Stokes system in

$$(0, T) \times \Omega, \quad \Omega = (0,1) \times \mathcal{T}^{N-1},$$

with the initial data

$$\varrho_0 \in L^\infty(\Omega), \quad \varrho_0 > 0, \quad \mathbf{u}_0 \in L^2(\Omega; \mathbb{R}^3).$$

Stability of 1D solutions - conclusion

Conclusion

Then

$$\int_{\Omega} \left[\frac{1}{2} \varrho |\mathbf{u} - \mathbf{V}|^2 + P(\varrho) - P'(R)(\varrho - R) - P(R) \right] (\tau, \cdot) \, dx$$

$$\leq c(T) \int_{\Omega} \left[\frac{1}{2} \varrho_0 |\mathbf{u}_0 - \mathbf{V}_0|^2 + P(\varrho_0) - P'(R_0)(\varrho_0 - R_0) - P(R_0) \right] \, dx$$

for a.a. $\tau \in (0, T)$,

$$P(\varrho) = \frac{a}{\gamma - 1} \varrho^{\gamma}.$$

Full Navier-Stokes-Fourier system

Navier-Stokes-Fourier system

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0,$$

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p(\varrho, \vartheta) = \operatorname{div}_x \mathbb{S}(\vartheta, \nabla_x \mathbf{u}) - \lambda \mathbf{u}$$

$$\partial_t(\varrho s(\varrho, \vartheta)) + \operatorname{div}_x(\varrho s(\varrho, \vartheta) \mathbf{u}) + \operatorname{div}_x \left(\frac{\mathbf{q}}{\vartheta} \right) = \sigma$$

$$\sigma = \frac{1}{\vartheta} \left(\mathbb{S}(\vartheta, \nabla_x \mathbf{u}) : \nabla_x \mathbf{u} - \frac{\mathbf{q} \cdot \nabla_x \vartheta}{\vartheta} \right),$$

Slip boundary conditions

$$\mathbf{u} \cdot \mathbf{n}|_{\partial\Omega} = 0, \quad [\mathbb{S}(\vartheta, \nabla_x \mathbf{u}) \cdot \mathbf{n}] \times \mathbf{n}|_{\partial\Omega} = 0$$

$$\mathbf{q}(\vartheta, \nabla_x \vartheta) \cdot \mathbf{n}|_{\partial\Omega} = 0$$

Constitutive relations - scaling

Pressure

$$p(\varrho, \vartheta) = p_M(\varrho, \vartheta) + p_R(\varrho, \vartheta), \quad p_M = \vartheta^{5/2} P\left(\frac{\varrho}{\vartheta^{3/2}}\right), \quad p_R(\varrho, \vartheta) = \frac{a}{3} \vartheta^4$$

Viscous stress

$$\mathbb{S}(\vartheta, \nabla_x \mathbf{u}) = \boxed{\nu} \left[\mu(\vartheta) \left(\nabla_x \mathbf{u} + \nabla_x^t \mathbf{u} - \frac{2}{3} \operatorname{div}_x \mathbf{u} \mathbb{I} \right) + \eta(\vartheta) \operatorname{div}_x \mathbf{u} \mathbb{I} \right]$$

Heat flux

$$\mathbf{q} = -\boxed{\omega} \kappa(\vartheta) \nabla_x \vartheta$$

Brinkman type “damping”

$$D = -\boxed{\lambda} \mathbf{u}$$

Target system

Full Euler system

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p_M(\varrho, \vartheta) = 0$$

$$\partial_t \left(\frac{1}{2} \varrho |\mathbf{u}|^2 + \varrho e_M(\varrho, \vartheta) \right)$$

$$+ \operatorname{div}_x \left[\left(\frac{1}{2} \varrho |\mathbf{u}|^2 + \varrho e_M(\varrho, \vartheta) \right) \mathbf{u} + p_M(\varrho, \vartheta) \mathbf{u} \right] = 0$$

Slip boundary conditions

$$\mathbf{u} \cdot \mathbf{n}|_{\partial\Omega} = 0$$

Dissipative solutions

Relative energy

$$\begin{aligned} & \mathcal{E}(\varrho, \vartheta, \mathbf{u} | r, \Theta, \mathbf{U}) \\ &= \int_{\Omega} \left[\frac{1}{2} \varrho |\mathbf{u} - \mathbf{U}|^2 + H_{\Theta}(\varrho, \vartheta) - \frac{\partial H_{\Theta}(r, \Theta)}{\partial \varrho} (\varrho - r) - H_{\Theta}(r, \Theta) \right] dx \end{aligned}$$

Relative energy inequality

$$\begin{aligned} & \left[\mathcal{E}(\varrho, \vartheta, \mathbf{u} | r, \Theta, \mathbf{U}) \right]_{t=0}^{t=\tau} \\ &+ \int_0^{\tau} \int_{\Omega} \frac{\Theta}{\vartheta} \left(\mathbb{S}(\vartheta, \nabla_x \mathbf{u}) : \nabla_x \mathbf{u} - \frac{\mathbf{q}(\vartheta, \nabla_x \vartheta) \cdot \nabla_x \vartheta}{\vartheta} \right) dx dt \\ &+ \lambda \int_0^{\tau} \int_{\Omega} |\mathbf{u}|^2 dx dt \leq \int_0^{\tau} \mathcal{R}(\varrho, \vartheta, \mathbf{u}, r, \Theta, \mathbf{U}) dt \\ & r, \Theta > 0, \mathbf{U} \cdot \mathbf{n}|_{\partial\Omega} = 0 \end{aligned}$$

Dissipative solutions - remainder

Remainder

$$\begin{aligned} \mathcal{R}(\varrho, \vartheta, \mathbf{u}, r, \Theta, \mathbf{U}) &= \int_{\Omega} \varrho (\mathbf{u} - \mathbf{U}) \cdot \nabla_x \mathbf{U} \cdot (\mathbf{U} - \mathbf{u}) \, dx \\ &+ \int_{\Omega} \left[\mathbb{S}(\vartheta, \nabla_x \mathbf{u}) : \nabla_x \mathbf{U} - \frac{\mathbf{q}(\vartheta, \nabla_x \vartheta)}{\vartheta} \cdot \nabla_x \Theta + \lambda \mathbf{u} \cdot \mathbf{U} \right] \, dx \\ &+ \int_{\Omega} \varrho \left(s(\varrho, \vartheta) - s(r, \Theta) \right) (\mathbf{U} - \mathbf{u}) \cdot \nabla_x \Theta \, dx \\ &+ \int_{\Omega} \varrho \left(\partial_t \mathbf{U} + \mathbf{U} \cdot \nabla_x \mathbf{U} \right) \cdot (\mathbf{U} - \mathbf{u}) \, dx - \int_{\Omega} p(\varrho, \vartheta) \operatorname{div}_x \mathbf{U} \, dx \\ &- \int_{\Omega} \left(\varrho \left(s(\varrho, \vartheta) - s(r, \Theta) \right) \partial_t \Theta + \varrho \left(s(\varrho, \vartheta) - s(r, \Theta) \right) \mathbf{U} \cdot \nabla_x \Theta \right) \, dx \\ &+ \int_{\Omega} \left(\left(1 - \frac{\varrho}{r} \right) \partial_t p(r, \Theta) - \frac{\varrho}{r} \mathbf{u} \cdot \nabla_x p(r, \Theta) \right) \, dx \end{aligned}$$

Vanishing dissipation limit

Theorem EF [2015]

Let $[\varrho_E, \vartheta_E, \mathbf{u}_E]$ be the classical solution of the Euler system in a time interval $(0, T)$, with the initial data $[\varrho_{0,E}, \vartheta_{0,E}, \mathbf{u}_{0,E}]$. Let $[\varrho, \vartheta, \mathbf{u}]$ be a weak (dissipative) solution of the Navier-Stokes-Fourier system, with the initial data $[\varrho_0, \vartheta_0, \mathbf{u}_0]$.

Then

$$\begin{aligned} & \mathcal{E} \left(\varrho, \vartheta, \mathbf{u} \mid \varrho_E, \vartheta_E, \mathbf{u}_E \right) (\tau) \\ & \leq c_1(T, \text{data}) \mathcal{E} \left(\varrho_0, \vartheta_0, \mathbf{u}_0 \mid \varrho_{0,E}, \vartheta_{0,E}, \mathbf{u}_{0,E} \right) \\ & + c_2(T, \text{data}) \max \left\{ a, \nu, \omega, \lambda, \frac{\nu}{\sqrt{a}}, \frac{\omega}{a}, \left(\frac{a}{\sqrt{\nu^3 \lambda}} \right)^{1/3} \right\} \end{aligned}$$

for a.a. $\tau \in (0, T)$.