

Symmetries of string compactifications and generalization of Freed-Witten anomaly

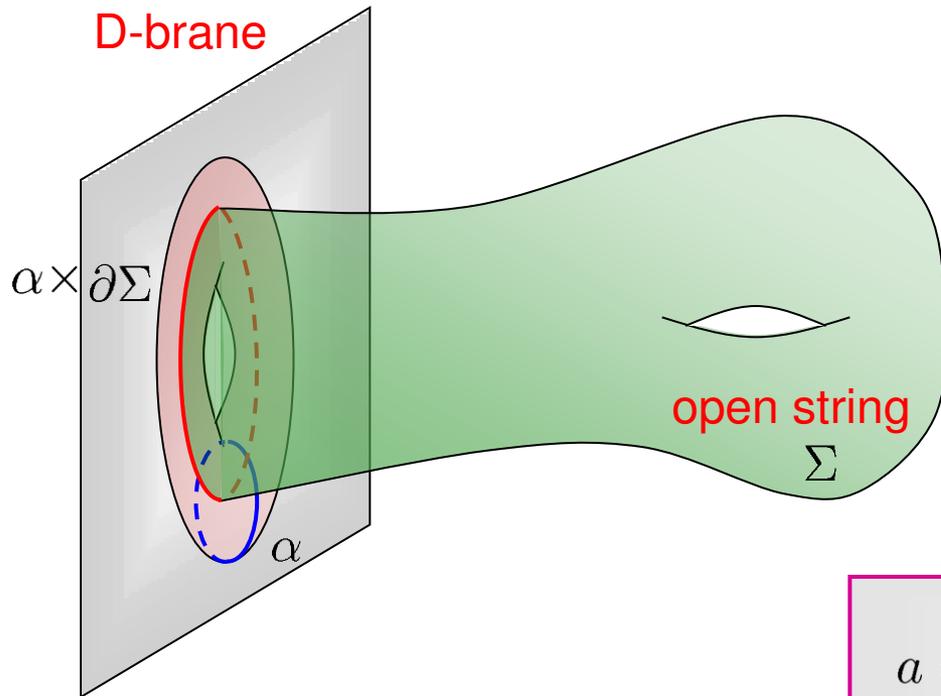
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S.A., S.Banerjee arXiv:1405.0291

Freed-Witten anomaly

Freed, Witten '99



worldsheet Dirac operator

U(1) gauge field on the brane

$$\text{pfaff}(D) \exp\left(2\pi i \oint_{\partial\Sigma} A\right)$$



going around a loop in the moduli space

$$(-1)^a \text{pfaff}(D)$$

$$a = \int_{\alpha \times \partial\Sigma} w_2(\text{D-brane})$$

Stiefel-Whitney class

path integral is well-defined

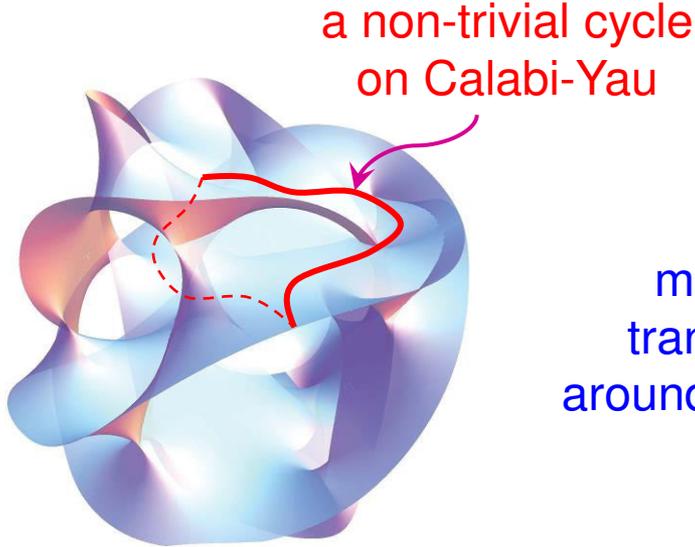


under variation of moduli periods of gauge field can get *half-integer* shifts

$$\oint_{\partial\Sigma} A \mapsto \oint_{\partial\Sigma} A + \frac{a}{2}$$

A is not an ordinary U(1) gauge field, but a Spin^c connection

Our result: a new anomaly



a non-trivial cycle
on Calabi-Yau

$$C = \int_{\gamma} A^{\text{even}} e^{-B}$$

gauge potentials of
type IIB string theory

RR-moduli of
the HM moduli space

monodromy
transformation
around large volume



large gauge transformation of
the B -field

$$\int_{\gamma^a} B \mapsto \int_{\gamma^a} B + \epsilon^a$$

symplectic transformation + anomaly

$$C \mapsto \rho(M_{\epsilon^a}) \cdot C + \text{half-integer terms}$$

axionic couplings of

D3-branes
(wrapping 4-cycles)

D5-branes
(wrapping CY)

$$C_a \mapsto \text{hom. terms} + A_{ab} \epsilon^b$$

$$C_0 \mapsto \text{hom. terms} + \frac{c_{2,a}}{8} \epsilon^a - \frac{1}{2} A_{ab} \epsilon^a \epsilon^b$$

second Chern
class of CY

half-integer matrix

$$A_{ab} = \frac{1}{2} \kappa_{abb} + \mathbb{Z}$$

Calabi-Yau compactifications

Type II string theory

compactification
on a Calabi-Yau



$\mathcal{N}=2$ supergravity in 4d
coupled to matter

- supergravity multiplet (metric & graviphoton)
- vector multiplets (gauge fields & scalars)
- hypermultiplets (only scalars)

The low energy effective action is completely
determined by the geometry of
moduli space

$$\mathcal{M}_{\text{VM}} \times \mathcal{M}_{\text{HM}}$$

In type IIB

\mathcal{M}_{VM} — special Kähler manifold parametrized by complex structure moduli

\mathcal{M}_{HM} — quaternion-Kähler manifold parametrized by

$$a = 1, \dots, h^{1,1}$$

$$\dim \mathcal{M}_{\text{HM}} = 4(h^{1,1} + 1)$$

$v^a = b^a + it^a$	complexified Kähler moduli
$c^0, c^a, \tilde{c}_a, \tilde{c}_0$	periods of RR gauge potentials
ψ	NS-axion (dual to the B-field)
τ_2	dilaton (string coupling $\tau_2 \sim g_s^{-1}$)

$$c^0 \equiv \tau_1$$



$$\tau = \tau_1 + i\tau_2$$

axio-dilaton

Classical symmetries

In the classical limit ($\alpha' \rightarrow 0$, $g_s \rightarrow 0$) the metric on \mathcal{M}_{HM} is given by the *c-map*. It has the following *continuous isometries*:

- **S-duality group $SL(2, \mathbb{R})$**

determined by the prepotential $F(X)$ describing Kähler moduli

$$\tau \mapsto \frac{a\tau + b}{c\tau + d} \quad t^a \mapsto t^a |c\tau + d| \quad \tilde{c}_a \mapsto \tilde{c}_a$$

$$\begin{pmatrix} c^a \\ b^a \end{pmatrix} \mapsto \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} c^a \\ b^a \end{pmatrix} \quad \begin{pmatrix} \tilde{c}_0 \\ \psi \end{pmatrix} \mapsto \begin{pmatrix} d & -c \\ -b & a \end{pmatrix} \begin{pmatrix} \tilde{c}_0 \\ \psi \end{pmatrix}$$

$$g = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

- **Heisenberg symmetry**

generated by shifts of NS-axion $k = \partial_\psi$ and RR-fields

$$h_{\eta^\Lambda, \tilde{\eta}_\Lambda} = \eta^\Lambda \partial_{c^\Lambda} + \tilde{\eta}^\Lambda \partial_{\tilde{c}^\Lambda} + \frac{1}{2} (\tilde{\eta}_\Lambda c^\Lambda - \eta^\Lambda \tilde{c}_\Lambda) \partial_\psi$$

$$[h_{\eta^\Lambda, 0}, h_{0, \tilde{\eta}_\Sigma}] = \delta_{\Sigma}^\Lambda k$$

$$\Lambda = (0, a) = 0, \dots, h^{1,1}$$

- **Shifts of the B-field**

$M_{\epsilon^a} : b^a \mapsto b^a + \epsilon^a$ supplemented by a change of RR-fields and NS-axion

Remarks

- **redundancy**

$$\begin{array}{ccc} \begin{array}{c} \curvearrowright \\ \text{Heisenberg} \\ \text{shift} \end{array} & H_{\eta^0, 0} \equiv e^{h_{\eta^0, 0}} = T^{\eta^0} & \begin{array}{c} \curvearrowleft \\ \text{SL}(2, \mathbb{R}) \\ \text{generator} \end{array} \end{array}$$

- **group structure**

$$\begin{array}{ccc} \begin{array}{c} \curvearrowright \\ \text{S-duality} \\ \text{subgroup} \end{array} & SL(2, \mathbb{R}) \times N(\mathbb{R}) & \begin{array}{c} \curvearrowleft \\ \text{nilpotent} \\ \text{subgroup} \end{array} \end{array}$$

Quantum symmetries

Quantum corrections break the continuous isometries to *discreet* subgroups

$$SL(2, \mathbb{R}) \times N(\mathbb{R}) \longrightarrow SL(2, \mathbb{Z}) \times N(\mathbb{Z})$$

But it is *not* sufficient to replace real transformation parameters by integers!

Amendments: (S.A., Persson, Pioline '10)

$$g \in SL(2, \mathbb{Z}) : \tilde{c}_a \mapsto \tilde{c}_a - c_{2,a} \varepsilon(g)$$

log of the multiplier system of Dedekind eta function

$$H_\Theta : \psi \mapsto \psi + \frac{1}{2} \langle \Theta, C \rangle + \sigma(\Theta)$$

quadratic refinement

$$\Theta = (\eta^\Lambda, \tilde{\eta}_\Lambda)$$

$$\langle \Theta, \Theta' \rangle = \tilde{\eta}_\Lambda \eta'^\Lambda - \eta^\Lambda \tilde{\eta}'_\Lambda$$

$$M_{\epsilon^a} : \psi \mapsto \psi + \frac{1}{6} \kappa_{abc} \epsilon^a c^b c^c + \kappa(M_{\epsilon^a})$$

character of symplectic group

Holomorphic prepotential:

$$F(X) = -\frac{\kappa_{abc} X^a X^b X^c}{6X^0} + \frac{1}{2} A_{\Lambda\Sigma} X^\Lambda X^\Sigma + \alpha' \text{-corr.}$$

$$A_{00} \in \mathbb{Z}$$

$$A_{0a} = \frac{c_{2,a}}{24} + \mathbb{Z}$$

$$A_{ab} = \frac{1}{2} \kappa_{abb} + \mathbb{Z}$$

affects mirror map

$$\tilde{c}_\Lambda^{\text{IIB}} = \tilde{c}_\Lambda^{\text{IIA}} - A_{\Lambda\Sigma} c^\Sigma$$

needs a correction

$$H_{\eta^0,0} \neq T^{\eta^0}$$

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character of symplectic group

Quadratic refinement $s(\Theta) = e^{2\pi i \sigma(\Theta)}$ of the intersection form on the charge lattice

$$s(\Theta_1 + \Theta_2) = (-1)^{\langle \Theta_1, \Theta_2 \rangle} s(\Theta_1) s(\Theta_2)$$

- multiplies D-instantons

$$s(\gamma) \Omega(\gamma) e^{-2\pi |Z_\gamma|/g_s - 2\pi i \langle \gamma, C \rangle}$$

crucial for consistency with wall-crossing

- needed to cancel the anomaly in the NS-axion shift:

$$H_{\Theta_1} H_{\Theta_2} = H_{\Theta_1 + \Theta_2} e^{\kappa \partial_\psi} \quad \text{where}$$

$$\kappa = \frac{1}{2} \langle \Theta_1, \Theta_2 \rangle + \sigma(\Theta_1) + \sigma(\Theta_2) - \sigma(\Theta_1 + \Theta_2)$$

= integer due to the quadratic refinement

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character of symplectic group

$\mathcal{M}_{\text{HM}} - S^1_\psi$ -bundle with curvature

$$F = \omega_C + \frac{\chi_C \gamma}{24} \omega_K$$

Kähler form on the torus of RR-fields

Kähler form for compl. Kähler moduli (1-loop effect)



ψ is not globally defined and the character allows to cancel fractional shifts due to monodromies in the Kähler moduli subspace

Since our group is abelian

$$\kappa(M_{\epsilon^a}) = \kappa_a \epsilon^a$$

Anomaly

Let us consider

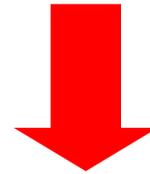
S-duality
generator

$$S^{-1} H_{\eta^a} S \cdot \begin{pmatrix} b^a \\ c^a \\ \tilde{c}_a \\ \tilde{c}_0 \\ \psi \end{pmatrix} = M_{\eta^a} \cdot \begin{pmatrix} b^a \\ c^a \\ \tilde{c}_a \\ \tilde{c}_0 \\ \psi \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ A_{ab} \eta^b \\ \eta^a \left(\frac{c_{2,a}}{8} - \frac{1}{2} A_{ab} \eta^b \right) \\ \eta^a \left(\frac{c_{2,a}}{24} - \kappa_a \right) \end{pmatrix} \text{ half-integer}$$

generate shifts of

$$\underbrace{c^a \quad b^a}$$

transform as a doublet
under $SL(2, \mathbb{Z})$



Anomaly
the group law fails
to hold

Resolution of the anomaly

The idea: the problem lies in the quadratic refinement

Defining equation

$$s(\Theta_1 + \Theta_2) = (-1)^{\langle \Theta_1, \Theta_2 \rangle} s(\Theta_1)s(\Theta_2) \quad \rightarrow$$

General solution

$$\sigma(\Theta) = \frac{1}{2} \eta^\Lambda \tilde{\eta}_\Lambda + \eta^\Lambda \phi_\Lambda - \tilde{\eta}_\Lambda \theta^\Lambda$$

characteristics
(generalized spin structure)

$s(\Theta)$ must be symplectic invariant

$$\rho = \begin{pmatrix} \mathcal{D} & \mathcal{C} \\ \mathcal{B} & \mathcal{A} \end{pmatrix} : \begin{pmatrix} \theta^\Lambda \\ \phi_\Lambda \end{pmatrix} \mapsto \rho \cdot \left[\begin{pmatrix} \theta^\Lambda \\ \phi_\Lambda \end{pmatrix} - \frac{1}{2} \begin{pmatrix} (\mathcal{A}^T \mathcal{C})_d \\ (\mathcal{D}^T \mathcal{B})_d \end{pmatrix} \right]$$

inhomogeneous transformation
of characteristics

inconsistent with $M_{\eta^a} = S^{-1} H_{\eta^a} S$

do not affect
characteristics

Observation: the characteristics and RR-fields always appear in combination

$$c^\Lambda - \theta^\Lambda \mapsto c^\Lambda$$

$$\tilde{c}_\Lambda - \phi_\Lambda \mapsto \tilde{c}_\Lambda$$

redefinition

The redefined fields transform *inhomogeneously* under symplectic group

$$\begin{pmatrix} 0 \\ 0 \\ A_{ab} \epsilon^b \\ \frac{c_{2,a}}{8} \epsilon^a - \frac{1}{2} A_{ab} \epsilon^a \epsilon^b \end{pmatrix}$$

inhomogeneous term for M_{ϵ^a}

All anomalies cancel!

provided $\kappa_a = \frac{c_{2,a}}{24}$

Coclusions

- We showed that the consistent implementation of discrete isometries on the HM moduli space of Calabi-Yau compactifications requires a modification of the *monodromy* transformations of the RR-fields by *inhomogeneous* terms determined by the second Chern class and intersection numbers.
- After this modification, it was shown that NS5-brane instantons derived using S-duality are also consistent with the Heisenberg and monodromy invariance.
- The anomalous transformation of the RR-fields provides a generalization of the Freed-Witten anomaly.

Can one derive the anomalous terms from the Pfaffian of the Dirac operator?